



Delhi Institute for Administrative Services  
India's Leading Institute for Civil Services Examination

ALL INDIA TEST SERIES CSE-2023

Candidate 's Information

1. NAME:- Aniket Hirde
2. UPSC ROLL NO:- 6701477
3. MOBILE NO:-
4. SUBJECT:- Paper 2 Full Length test
5. DATE:- 27-08-23

Dias Roll No: 230001

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Q.NO	MARKS
1.	24
2.	
3.	28
4.	
5.	29½
6.	32
7.	30½
8.	

*Very Good*  
TOTAL MARKS 144

250

EXAMINER SIGNATURE

INVIGILATOR SIGNATURE

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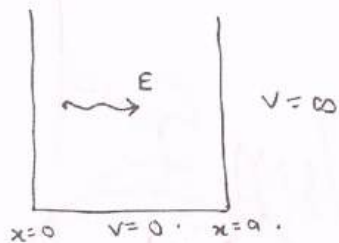
# UPSC

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उम्मीदवारों को इस हाशिए में नहीं लिखना चाहिए  
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Q1 (a)

Bohr's correspondence principle states that for a particle, energy states are quantized i.e. he gave the concept of stationary states



For electron in a box,

$$-\frac{\hbar^2}{2m} \frac{d^2 \psi}{dx^2} = E \psi$$

$$\Rightarrow \frac{d^2 \psi}{dx^2} = -k^2 \psi \quad \left\} \quad k = \frac{\sqrt{2mE}}{\hbar}$$

Fig 1.  $e^-$  in box

$$\Rightarrow \psi = A \sin kx + B \cos kx$$

Boundary condition:  $\psi|_{x=0} = 0 \Rightarrow B = 0$

$$\Rightarrow \psi = A \sin kx = 0 \text{ at } x = a$$

$\Rightarrow ka = n\pi$ . Substituting value of k,

$$E = \frac{n^2 \pi^2 \hbar^2}{2ma^2}$$

$$\Rightarrow \boxed{E = n^2 (\text{constant})}$$

Hence, the states formed are stationary in line with Bohr's model.

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Q1  
(b)

De-Broglie's wave-Particle Duality states that every moving particles behaves like a wave with wavelength  $(\lambda) = h/p$ .

(i)  $V = 10V$  i.e. non relativistic regime as  $E < m_0c^2$ .

$$\Rightarrow E = \frac{p^2}{2m} \Rightarrow p = \sqrt{2mE} \Rightarrow \lambda = \frac{h}{\sqrt{2mE}}$$

Substituting values,  $\lambda = 3.88 \text{ \AA}$

(ii)  $V = 1 \text{ MeV}$  as  $E > m_0c^2$ , relativistic.

$$\Rightarrow E^2 = p^2c^2 + (m_0c^2)^2 \Rightarrow (pc)^2 = 1 - (0.51)^2$$

$$\Rightarrow pc = 0.86 \text{ MeV}$$

$$\Rightarrow \lambda = \frac{hc}{pc} \Rightarrow \text{wavelength} = 1.44 \times 10^{-12} \text{ m}$$

As seen, in relativistic regime, momentum increases which leads to drop in wavelength.

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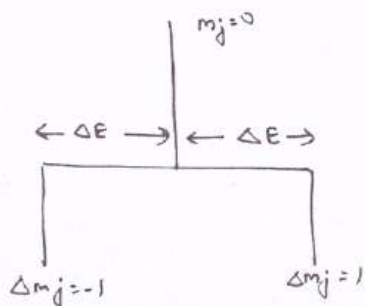
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Q1  
(c)

Zeeman splitting is splitting of  $j^{\text{th}}$  energy level into  $(2j+1)$  sublevels in the presence of weak magnetic field.



> In normal Zeeman,  
 $S = 0 \Rightarrow J = L$   
 $\Rightarrow g_J = 1$

> Energy of splitting

$$E = -\mu_J \cdot B$$

$$\Rightarrow \Delta E = \frac{e}{2m} g_J m_J \hbar B$$

$$\Rightarrow \Delta E = \mu_B B \quad \text{Also } \Delta E = d \left( \frac{hc}{\lambda} \right) = \frac{hc \Delta \lambda}{\lambda^2}$$

$$\Rightarrow \Delta \lambda = \frac{(\mu_B B) (\lambda^2)}{hc} = \underline{0.33 \text{ \AA}}$$

$\therefore$  Wavelengths of components are  $\lambda, \lambda + \Delta \lambda$

&  $\lambda - \Delta \lambda$  i.e.  $\boxed{4226.4 \text{ \AA}}$ ,  $\boxed{4226.73 \text{ \AA}}$

&  $\boxed{4227.06 \text{ \AA}}$ . The separation

between them would be  $\boxed{0.33 \text{ \AA}}$ .

Zeeman effect is used  
in laser cooling application.

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Q.1  
(d)

Postulates of quantum mechanics state that every measurable quantity can be represented by a wave function

(i) Expectation value of <sup>weighted by probability</sup> any physical quantity is the value that would be obtained if it is measured multiple times.

> If the wave function is  $\psi$ , expectation value of any quantity  $A$  is  $\langle \psi | A | \psi \rangle$ .

By  $\langle A \rangle$ , we can know,

(i) Shape of wave function

(ii) Error in measurement if  $\langle A^2 \rangle$  known.

(ii) Given:-  $\psi = ax$  for  $0 < x < 1$

We know, probability =  $\psi \psi^*$ . As

particle will be found somewhere,

$$\langle \psi | \psi^* \rangle = 1 \Rightarrow \int_{-\infty}^0 \psi^* \psi + \int_0^1 \psi \psi^* + \int_1^{\infty} \psi^* \psi = 1$$

$$\Rightarrow \boxed{\int \psi \psi^* = 1} \quad \therefore \boxed{\text{Prob}(x=0 \text{ to } x=1) = 1}$$

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Q1(e)

Potential energy is the energy of electron due to Coulombic interaction with proton.

$$\text{We know, } V = \frac{-e^2}{4\pi\epsilon_0 r} \Rightarrow \langle V \rangle = \frac{-e^2}{4\pi\epsilon_0} \left\langle \frac{1}{r} \right\rangle$$

$$\Rightarrow \left\langle \frac{1}{r} \right\rangle = \int_0^\infty \frac{1}{r} \cdot \frac{1}{\pi a_0^3} e^{-2r/a_0} \cdot 4\pi r^2 dr$$

$$\Rightarrow \left\langle \frac{1}{r} \right\rangle = \frac{4}{a_0^3} \int_0^\infty r e^{-2r/a_0} dr \quad \left. \begin{array}{l} \text{Let } \frac{2r}{a_0} = x \\ dr = \frac{a_0}{2} dx \end{array} \right\}$$

$$\Rightarrow \left\langle \frac{1}{r} \right\rangle = \frac{4}{a_0^3} \cdot \left(\frac{a_0}{2}\right)^2 \int_0^\infty x e^{-x} dx$$

$$\Rightarrow \left\langle \frac{1}{r} \right\rangle = \frac{1}{a_0} = (0.529 \text{ \AA})^{-1}$$

$$\Rightarrow \langle V \rangle = \frac{-e^2}{4\pi\epsilon_0} \times \frac{me^2}{\hbar^2} = \frac{-me^4}{4\epsilon_0\hbar^2}$$

Given.

Substituting values,  $\langle V \rangle = \boxed{-27.2 \text{ eV}}$

As seen, the potential energy is twice of total energy.

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Q3

(a)

Particle in potential well is characterized by potential  $\propto x^2$  where  $x$  is the distance. Hence, Schrodinger Wave Eq. is

$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} + \left(\frac{1}{2} kx^2\right) \psi = E\psi.$$

But we define 2 operators as :-

$$a_+ = \text{Raising} = \frac{1}{\sqrt{2m\hbar\omega}} [m\omega x - iP_x]$$

$$a_- = \text{Lowering} = \frac{1}{\sqrt{2m\hbar\omega}} [m\omega x + iP_x]$$

Let The operator operate as :-

$$a_+|n\rangle = \sqrt{n+1}|n+1\rangle \quad \& \quad a_-|n\rangle = \sqrt{n}|n-1\rangle$$

for state  $n$ . Let the ground state wave function be  $\psi_0$ .

$$\Rightarrow a_-|\psi_0\rangle = 0.$$

$$\Rightarrow m\omega x \psi_0 + i(-i\hbar \frac{\partial \psi_0}{\partial x}) = 0.$$

$$\Rightarrow m\omega x \psi_0 = -\hbar \frac{\partial \psi_0}{\partial x}$$

$$\Rightarrow \int \frac{\partial \psi_0}{\psi_0} = \int \frac{m\omega}{-\hbar} x dx \Rightarrow \underline{\underline{\psi_0 = Ae^{-\frac{m\omega}{2\hbar} x^2}}}}$$

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We know,  $\langle \Psi_0 | \Psi_0 \rangle = 1$

$$\Rightarrow \int_{-\infty}^{\infty} |A|^2 e^{-\frac{m\omega}{\hbar} x^2} dx = 1 \text{ . on solving ,}$$

$$\Rightarrow |A|^2 = \left( \frac{m\omega}{\pi\hbar} \right)^{1/4}$$

$$\Rightarrow \Psi_0 = \left( \frac{m\omega}{\pi\hbar} \right)^{1/4} e^{-\frac{m\omega}{2\hbar} x^2}$$

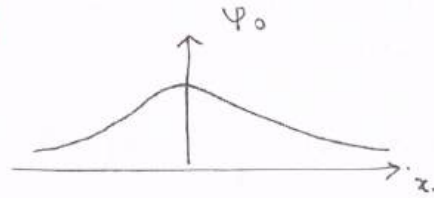


Fig 1.  $\Psi_0$  state

Also,  $a_+ a_- = \frac{\hbar}{2m\omega} - \frac{1}{2}$  . We know,  $H\Psi = E\Psi$

$$\Rightarrow \frac{\hbar\omega}{2} |\Psi_0\rangle = a_+ a_- + \frac{1}{2} |\Psi_0\rangle$$

$$\Rightarrow \frac{E}{\hbar\omega} \Psi_0 = \frac{\Psi_0}{2} \Rightarrow E = \frac{\hbar\omega}{2}$$

The quantum oscillator differs from classical as  $E_{\min} \neq 0$ , energy levels are quantized & probability of finding particle outside of max. amplitude  $\sim 16\%$ .

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Q3  
(b)

Particle in a finite well can be modeled using Schrodinger's wave Equation which gives symmetric & antisymmetric solutions -

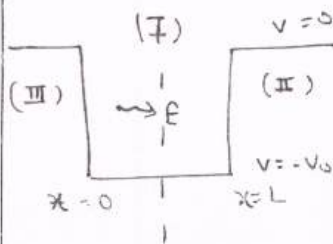


Fig 1. Particle in finite well

> We know, for region (I)

$$-\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} + V\psi = E\psi$$

$$\Rightarrow \psi_I = A \sin k_1 x + B \cos k_1 x$$

$$\text{Where } k_1 = \frac{\sqrt{2m(V_0 - E)}}{\hbar}$$

$$\psi_{II} = C e^{-k_2 x} \quad ; \quad k_2 = \frac{\sqrt{2mE}}{\hbar}$$

Boundary conditions :-  $\psi_I = \psi_{II} \Big|_{x=L}$

&  $\frac{\partial \psi_I}{\partial x} = \frac{\partial \psi_{II}}{\partial x} \Big|_{x=L}$  .. Here

2 cases arise :

① Anti-symmetric solutions :-  $\psi_I = A \sin k_1 x$

$$\Rightarrow A \sin k_1 L = C e^{-k_2 L}$$

$$-k_1 A \cos k_1 L = k_2 C e^{-k_2 L}$$

$$\Rightarrow x \tan\left(\frac{\pi}{2} + x\right) = y$$

where,  $x = k_1 L$

$y = k_2 L$

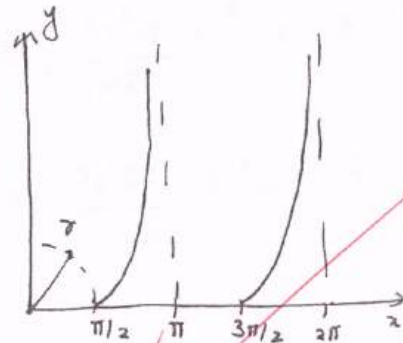


Fig 2. Anti-symmetric

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$\Rightarrow$  If  $x^2 + y^2 = \gamma^2$ , For  $\gamma < (n-1/2)\pi$ ,  
 We have,  $(n-1)$  Anti-symmetric states.

② Symmetric States:-  $\Psi_{II} = B \cos k_1 x$ .

$\Rightarrow B \cos k_1 L = C e^{-k_2 L}$

$B k_1 \sin k_1 L = C k_2 e^{-k_2 L}$

$\Rightarrow x \tan x = y$

where  $x = k_1 L$

$y = k_2 L$

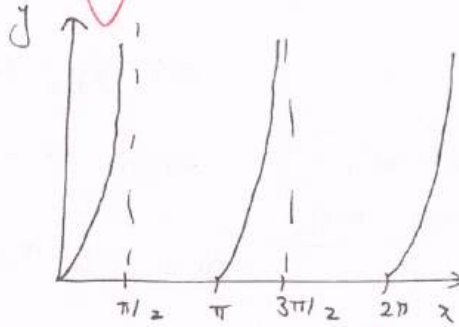


Fig 3. Symmetric states

$\Rightarrow$  For  $x^2 + y^2 = \gamma^2$

If  $\gamma < (n-1/2)\pi$ ,

$n$  symmetric bound-states.

Now,  $x^2 + y^2 = \frac{2m V_0}{\hbar^2} \cdot L^2 = \frac{2m (V_0 L^2)}{\hbar^2}$

Now if  $V_0 L^2 = \frac{\pi^2 \hbar^2}{8m}$ ,  $\gamma^2 = \frac{\pi^2}{4} \Rightarrow \gamma = \frac{\pi}{2}$

if  $V_0 L^2 = \frac{4\pi^2 \hbar^2}{2m}$ ,  $\gamma^2 = 4\pi^2 \Rightarrow \gamma = 2\pi$



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$\therefore$  For  $\frac{\pi}{2} < \alpha \leq 2\pi$ , we will have,

2 antisymmetric  $\oplus$  3 symmetric state

$\Rightarrow$  Total no. of bound states = 5.

Particle in finite well  
is used to model deuteron to  
estimate range of nuclear force.

$$r^2 = \frac{mV_0L^2}{2\hbar^2}$$

$$\frac{10}{25}$$

$$\frac{\pi}{4} < \alpha \leq \pi$$

only 3 solns.

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Q3(c)

Raman Spectra is observed when scattered light gives frequency both greater & less than incident light.

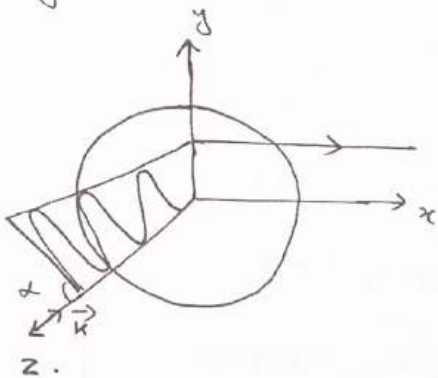


Fig. Light incident on molecule

> Consider a light with plane of polarization at  $\angle \alpha$  with z-axis falling on spherical molecule like  $\text{CH}_4$ .

> As the molecule is symmetric, its sphere of polarizability will expand equally when light falls on it.

> Its <sup>bonds</sup> will vibrate & in plane of incident light.

> As light emitted is  $\perp$  to plane of vibration, the plane would always be xy plane irrespective of  $\angle \alpha$ .

> However, not all molecules are spherically symmetric. In such cases,

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polarization would be maximum along minor axis of ellipsoid. & the light would be polarized to different extents.

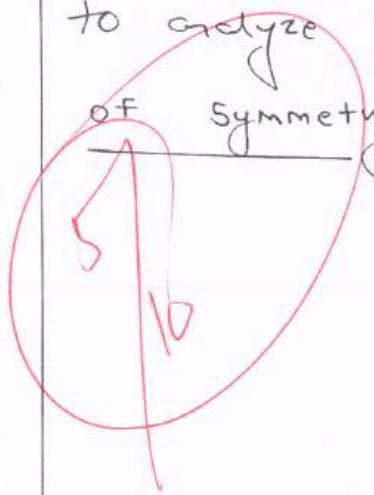
> Hence, light will always lie in xy plane (to different extent) & will be plane polarized.

*Degree of polarizability*

$$r = \frac{I_{\perp}}{I_{\parallel}}$$

Raman spectra is used

to analyze molecules without centre of symmetry.



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Q5(a)

Nuclear properties are the properties of nucleus i.e. a region where protons & neutrons exist.

① Volume  $\propto A$  :- Volume of nucleus is proportional to its mass no. (A). This is because, radius of nucleus  $R = R_0 A^{1/3}$   
 $\Rightarrow \text{Volume} = \frac{4}{3}\pi R^3 = \left(\frac{4}{3}\pi R_0^3\right) A$

② Constant density :- As mass of nucleus is  $\approx A$  a.m.u & volume  $\propto A$ ,  
$$\rho = \frac{A \text{ (a.m.u)}}{\frac{4}{3}\pi R_0^3 A} \approx \underline{\underline{10^{17} \text{ kg/m}^3}}$$
 i.e.  
density of nucleus across atoms is constant.

Numerical For  $^{12}\text{C}$ ,  $A = 12$

$$\Rightarrow R = R_0 A^{1/3} = \underline{\underline{2.75 \text{ fm}}}$$

$$\frac{\text{Atomic radius}}{\text{Nuclear radius}} = \frac{0.529 \times 10^{-10}}{2.75 \times 10^{-15}} = \underline{\underline{1.9 \times 10^4}}$$

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∴ Atomic radius of carbon is  $\sim 10^4$  times that of nuclear radius.

The inference that can be drawn is that nuclear <sup>force</sup> ~~range~~ is very short range i.e.  $10^{-15}$  m. However, coulombic force i.e.  $e-p$  force has longer range.



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Q.5  
(b)

The number of neutrons v/s number of protons plot can help estimate relative Stability of nucleus.

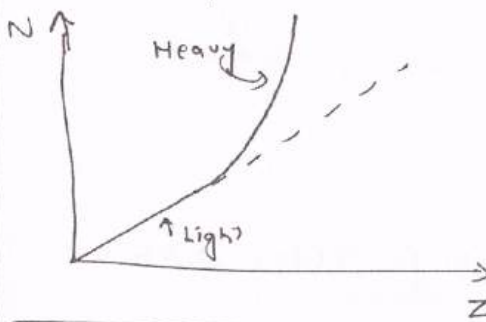


Fig 1. N v/s Z plot

(i) Inference

① For lighter nuclei,

$$Z < 20, N = Z.$$

② For heavier nuclei,  $N > Z$ .

because Coulombic repulsion increases more rapidly v/s nuclear attraction  $\Rightarrow$  More neutrons needed to balance repulsion.

(ii) Between  ${}^7_3\text{Li}$  &  ${}^8_3\text{Li}$ ,  ${}^7_3\text{Li}$  would be more stable as N & Z are closer.

For  ${}^8_3\text{Li}$ , nuclear forces will dominate the Coulombic force leading to collapse of atom.

This Z to N conversion gives the asymmetry term in Weizsacker-Bethe formula.



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Q5  
(c)

Cell is a fundamental unit of structure replicating whom along translation vector gives the whole crystal. Depending on placement of atoms, we have 3 types:-

## ① Simple Cubic

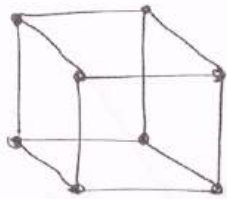


Fig 1. Simple Cubic

Here,  $2r = a$

$$\Rightarrow \# \text{atoms / cell} = 1$$

$$\Rightarrow \text{Packing fraction} = \frac{(1) \left( \frac{4}{3} \pi r^3 \right)}{a^3}$$

$$= 52\%$$

Example:- Polonium

## ② Body-Centered Cubic

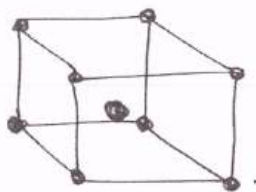


Fig 2. BCC cell

Here, atoms touch along body diagonal

$$\Rightarrow \sqrt{3}a = 4r$$

$$\Rightarrow \# \text{atoms / cell} = 2$$

$$\Rightarrow \text{Packing fraction} = \frac{2 \times \left( \frac{4}{3} \pi r^3 \right)}{a^3}$$

$$= 68\%$$

Example:- Na, K

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## ③ Face-centered cubic

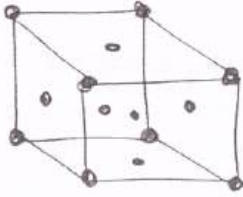


Fig 3. FCC cell

> Atoms touch along face diagonal  $\Rightarrow 4r = \sqrt{2}a$ .

$$\Rightarrow \# \text{ atoms / cell} = \boxed{4}$$

$$\Rightarrow P.f = \frac{4 \left( \frac{4}{3} \pi r^3 \right)}{a^3} = \boxed{74.1\%}$$

Example :- Cu, Ag.

Simple Cubic & BCC are called loose packed structure due to low packing fraction while FCC is closed packed structure.

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Q.5  
(cd)

Bound-state is a state for a particle when it is confined to particular region due to some potential.

> n-n scattering establishes that both di-proton & di-neutron do not exist.

> Can be explained via exchange force model :-

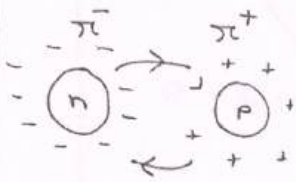


Fig 1. Exchange forces

> Particles like neutron & proton are attracted due to exchange of  $\pi^+$  mesons. This leads to bound-state.

> In case of neutron-neutron, no particles will be exchanged  $\Rightarrow$  No force.

> To account for it, iso-spin has been introduced with  $T_z = 1/2$  for proton &  $T_z = -1/2$  for neutron

This shows that nuclear force follows iso-spin symmetry.

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Q.5(e)

Gellman - Newmann proposed quark model of Hadrons stating that each hadron can be written as function of quarks.

The quarks & properties are:-

Quark	Charge	B	Spin	$T_z$	Strangeness
u	$\frac{2}{3}e$	$\frac{1}{3}$	$\frac{1}{2}$	$\frac{1}{2}$	0
d	$-\frac{1}{3}e$	$\frac{1}{3}$	$\frac{1}{2}$	$-\frac{1}{2}$	0
s	$-\frac{1}{3}e$	$\frac{1}{3}$	$\frac{1}{2}$	0	-1

(i)  $\pi^0$  :-  $Q=0, B=0, T_z=0, \text{Strangeness}=0$

$$\Rightarrow \pi^0 = \frac{1}{\sqrt{2}} (u\bar{u} - d\bar{d})$$

(ii)  $K^+$  :-  $Q=1, B=0, T_z=\frac{1}{2}, \text{Strangeness}=1$

$$\Rightarrow K^+ = u\bar{s}$$

(iii)  $\Delta^{++}$  :-  $Q=2, B=1, T_z=3/2, \text{Strangeness}=0$

$$\Rightarrow \Delta^{++} = uuu$$

(iii)  $\Sigma^0$  :-  $Q=0, B=1, T_z=0, \text{Strangeness}=-1$

$$\Rightarrow \Sigma^0 = uds$$

This model is further used to study decay of strange particles.

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Q 6  
(a)  
(i)

Bethe - Weizsacker gave semi-empirical mass  
Formula to estimate mass & binding  
energy of atoms as:-

$$M(Z, A) = Zm_p + (A-Z)m_n - \underbrace{a_1 A}_{E_v} + \underbrace{a_2 A^{2/3}}_{E_s} \\ + \underbrace{a_3 \frac{Z(Z-1)}{A^{1/3}}}_{E_c} + \underbrace{a_4 \frac{(A-2Z)^2}{A}}_{E_a} \pm \underbrace{a_5 A^{-3/4}}_{E_p}$$

① Volume energy ( $E_v$ ): Accounts of nuclear force  
& is  $\propto V \Rightarrow E_v \propto \frac{4}{3}\pi(R_0 A^{1/3})^3$   
 $\Rightarrow E_v = -a_1 A$  } As it is attractive.

② Surface energy ( $E_s$ ): Nucleons at surface interact  
with lesser no. of nucleons  $\Rightarrow E$  decreases  
 $\therefore E_s \propto \text{surface area} \Rightarrow E_s = a_2 A^{2/3}$

③ Coulombic energy ( $E_c$ ): Accounts for proton -  
proton interaction  $E_c = \frac{ZcZe^2}{4\pi\epsilon_0 \langle r_{ij} \rangle}$   
 $\Rightarrow E_c = a_3 \frac{Z(Z-1)}{A^{1/3}}$

④ Asymmetry energy ( $E_a$ ): At higher mass no.,  
no. of neutrons  $> Z$  to account for.

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increased repulsion. Also some protons convert into neutrons. For every  $\left(\frac{N-Z}{2}\right)$  conversions, we have

$$\Delta E = \left(\frac{N-Z}{2}\right)^2 (a)$$

$$\Rightarrow E_a = a_4 \frac{(N-Z)^2}{A}$$

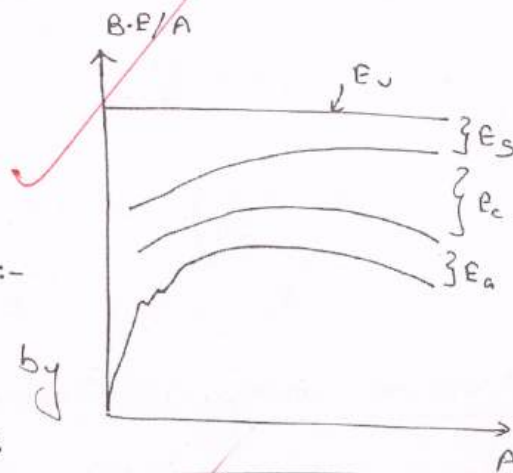


Fig 1. BE/A vs A

### (5) Pairing energy ( $E_p$ ):-

The stability shown by nucleus is N-Z as  
 even-even > even-odd >  
 odd-even > odd-odd

Hence  $E_p = +$  for ~~odd~~ even-even  
 $-$  for odd-odd nucleus  
 $0$  for odd A nucleus.

(ii) Exchange of pions leads to attractive force which is responsible for nuclear force.

> From Heisenberg's uncertainty principle, we know,  $\Delta E \cdot \Delta t \geq \hbar/a$ .



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$$\text{Now } \Delta E \sim E = m_{\pi}c^2, \quad \Delta t = \frac{\text{Range}}{c} = \frac{R}{c}.$$

$$\Rightarrow (m_{\pi}c^2) \cdot \frac{R}{c} \geq \frac{\hbar}{2}$$

$$\Rightarrow R \geq \frac{\hbar}{2m_{\pi}c} \quad \text{Substituting values,}$$

$$\text{We get range} = \boxed{0.7 \text{ fm}}$$

$\therefore$  Range of nuclear force is

$$\boxed{\sim 1 \text{ fm}}$$

nucleus.

which is also the size of



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Q. 6  
(b)

Shell-model analyzes one nucleon in potential of  $(A-1)$  nucleons. One of this type of potential is harmonic potential.

$$\Rightarrow \frac{-\hbar^2}{2m} \nabla^2 \psi + \left(\frac{1}{2} k r^2\right) \psi = E \psi. \text{ On solving,}$$

we get,  $E = (N + 3/2) \hbar \omega$

where  $N = \underline{2n + l - 2}$ . Here  $N = 0, 1, 2, \dots$   
 $l = 0, 1, 2, \dots$   
 $n = 0, 1, 2, \dots$

Solving for values of N :-

<u>N</u>	<u>n, l</u>	<u>State</u>	<u># Nucleons</u>
0	$n=1, l=0$	1s	2
1	$n=1, l=1$	1p	6
2	$n=1, l=2$	1d	10
	$n=2, l=0$	2s	2
3	$n=2, l=1$	2p	6
	$n=1, l=3$	1f	14
	⋮		



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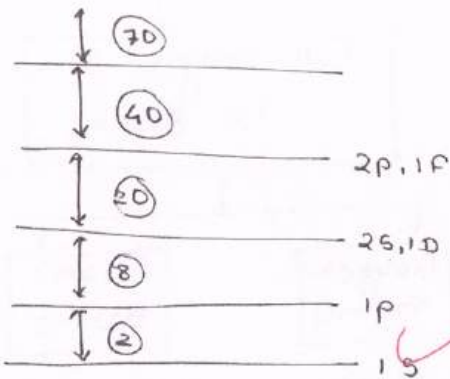
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As seen, for  $N = 1$ , we have 6 degenerate states, for  $N = 2$  we have 12 & so on. In general, for  $N$ , we have  $(N+1) \times (N+2)$  degenerate states for different values of  $n$  &  $l$ . Hence, there is



large degeneracy for particular  $N$ .

> The Shell enclosures & magic nos. predicted are 2, 8, 20, 40, 70

Fig 1. Magic numbers predicted

... → As seen they diverge from 2, 8, 20, 28, 50... for higher magic numbers

Mayer in 1949 introduced a non-central potential in the form of spin-orbit which led to splitting of  $I^{\text{th}}$  level to match magic numbers.

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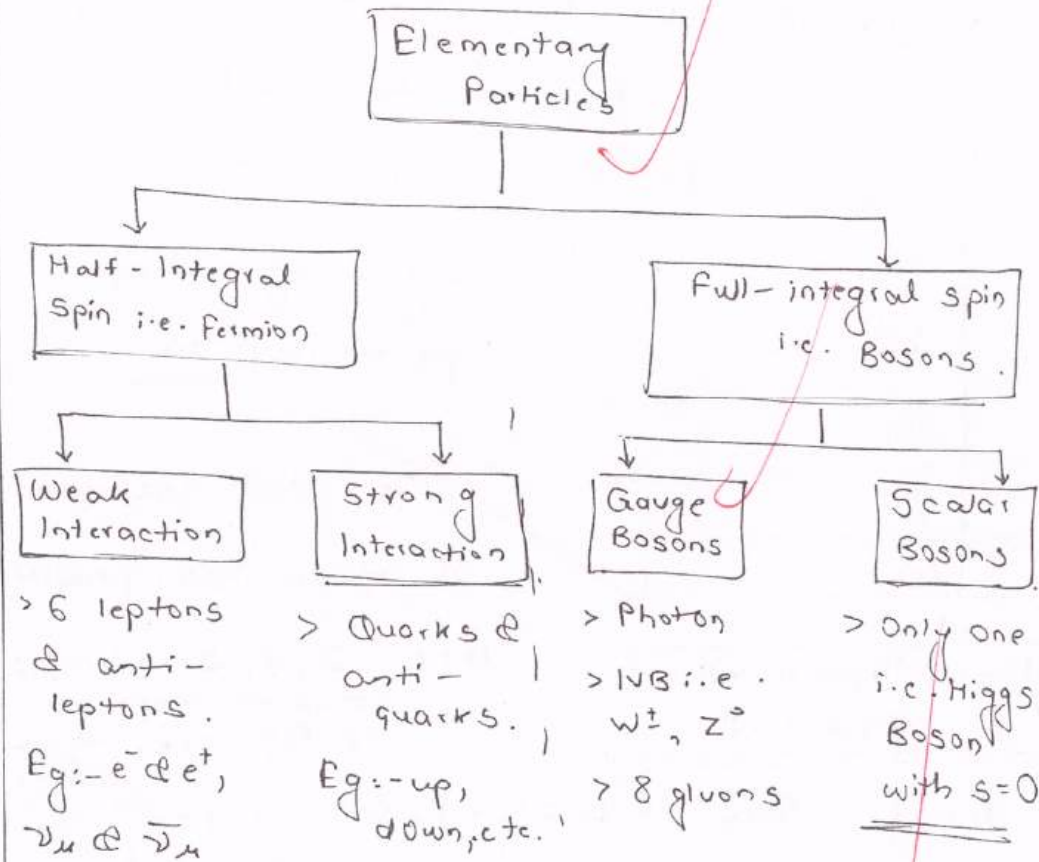
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Q.6  
(c)

Elementary particles are sub-atomic particles which cannot be divided further.



## Other types of classification

① Strangeness (Gellman - Nishijima).

Strangeness = 0 . Eg:- Leptons,  $\pi$ , etc.

Strangeness  $\neq 0$  Eg:-  $K^+$ ,  $\Sigma^0$ , etc. i.e.

Strange particles -

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② MASS :- Leptons ( $e^-$ ,  $\nu_e$ , etc.) < Mesons ( $\pi, K, \rho$ , etc.) < Baryons ( $n, p, \Lambda, \Sigma$ , etc.)

These elementary particles interact through various ways as:-

Interaction	$\tau$	Range	Particle	Exchange Particle	Strength
Strong	$10^{-23}$ s	$10^{-15}$ m	Hadrons	Gluons	1
Electro-magnetic	$< 10^{-14}$ s	$\infty$	Charged particles	$\gamma$	$10^{-4}$
Weak	$> 10^{-22}$ s	$10^{-17}$ m	Leptons	$W^\pm, Z^0$	$10^{-14}$
Gravity	-	$\infty$	'Particle' with mass	Graviton	$10^{-40}$

The properties of these interactions are:-

- ① Strong :- Short range, attractive, charge, mass independent, tensor.
- ② EM :- Both attractive, repulsive.
- ③ Weak :- Shortest range, frame variant, seen in strange particle decay.

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④ Gravity:- Always attractive but weakest.

Although standard model of physics assumes these to be different interactions, Grand Unification Theory states that they are manifestation of same force & would unify at  $10^{16}$  GeV.

Till now, electro-weak unification has been achieved at  $10^2$  GeV. Detection of Higgs Boson in 2012 has given new impetus to GUT.



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Q.7  
(a)  
(i)

Binding energy is the energy released when constituent nucleons come together in a nucleus.

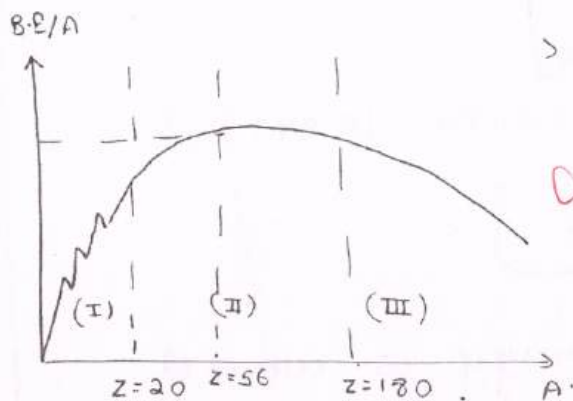


Fig. B.E/A vs A

> B.E vs A is an inverted parabola.

> Every nucleus tries to achieve more stability by increasing its Binding energy (B.E)

> If  $Z < 20$ , atoms are lighter  $\rightarrow$  Combine to increase B.E  $\Rightarrow$  Fusion. Eg:-  $^1_1\text{H}$

> If  $Z > 180$ , atoms are very heavy  $\rightarrow$  Break into constituent to decrease B.E.

$\Rightarrow$  fission. Eg:-  $^{235}_{92}\text{U}$ .

> Region (II) :- Nucleus are usually stable & don't show ~~radio~~ nuclear fission/fusion

Eg:-  $^{56}_{26}\text{Fe}$  has highest B.E & is very stable.

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(ii)

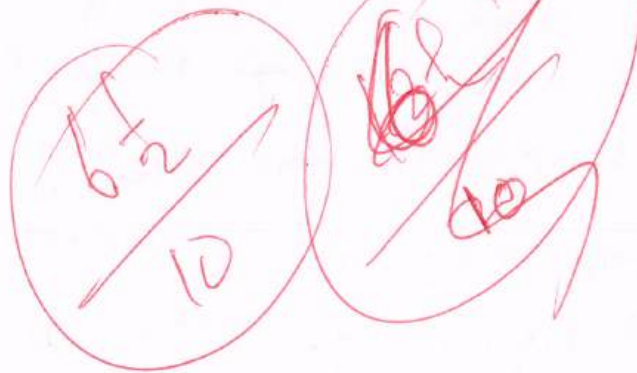
The energy of nuclear reaction can be given by the  $Q$ -value. We know,

$$Q = (m_i - m_f) c^2.$$

$$\Rightarrow Q = (12 + 4.002603 - 15.999) c^2$$

$$\Rightarrow \boxed{Q = 3.38 \text{ MeV}}$$

As energy is released in each reaction, eventually star runs out of fuel to form white dwarf, neutron star or black hole.



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Q.7  
 (b)

Josephson effect is tunnelling of Cooper pairs via thin insulator ( $\sim 10\text{\AA}$ ) sandwiched between two identical superconductors

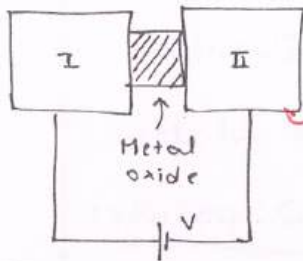


Fig. 1. Josephson junction

Here, 2 cases obtained:-

(A) DC Josephson i.e.  $V=0$

Even in this case, Cooper pairs tunnel through

(B) AC Josephson i.e.  $V=V_0$

We know, energy of Cooper pairs =  $2eV_0$

Now  $2eV_0 = h\omega$

$$\Rightarrow \omega = \frac{2eV_0}{\hbar} \therefore \text{If}$$

phase difference at  $t=0$  is  $\phi_0$ ,

We have current across

Josephson as 
$$J = J_0 \sin \left( \phi_0 \pm \frac{eV_0}{\hbar} t \right)$$

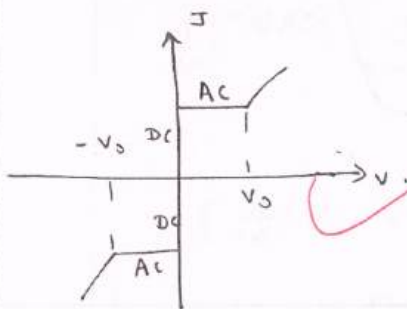


Fig. 2. AC & DC Josephson.

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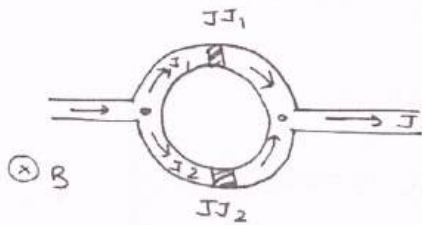
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This effect is further used to construct SQUID (Superconducting Quantum Interference Device) which are 2 Josephson connected in parallel.



> Magnetic Quantum Interference states that when 2 parallel

Fig 3. SQUID

Josephson junctions are

kept in constant magnetic field, the super-current interferes leading to maxima & minima.

> Experimentally, we know,  $\Phi = s \left( \frac{h}{2e} \right)$

$\Rightarrow 2\pi s = \frac{2e\Phi}{h}$ , we split this as :-

$$J_1 = J_0 \sin \left( \delta_0 - \frac{e\Phi}{h} \right) \quad \& \quad J_2 = J_0 \sin \left( \delta_0 + \frac{e\Phi}{h} \right)$$

$$\Rightarrow J = J_1 + J_2 = \underline{\underline{2J_0 \cos \left( \frac{e\Phi}{h} \right) \sin(\delta_0)}}$$



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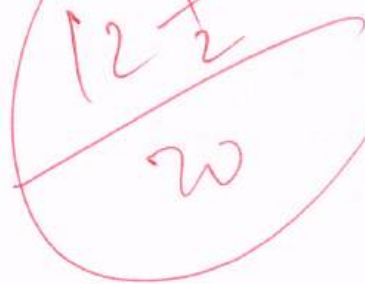
⇒ Every time  $\frac{e\Phi}{h} = \pi$ , we get a maxima

> We measure this maxima using galvanometer → find  $\Phi$ . If we know area, we can find magnetic field (B).

## Applications of JJ & SQUID

① Josephson junction used to measure extremely small voltages i.e.  $\mu V$  as  $V \propto \nu(Hz)$  which can be easily measured

② SQUID is used to measure extremely small magnetic fields like magnetic field at centre of atom due to rotating electron.



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उम्मीदवारों को इस हिसाब में नहीं लिखना चाहिए  
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Q 7  
(c)

Junction Field Effect Transistor (JFET) is a type of transistor in which conduction happens only via one type of carrier i.e. electron or hole.

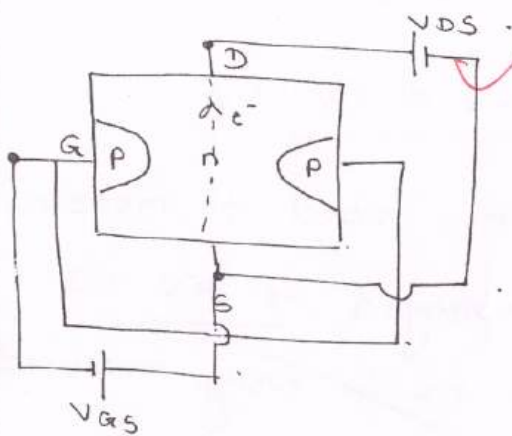


Fig. 1. n-channel JFET

> In n-channel JFET, we have 2 regions doped with p-type while the channel is n-type.

> It conducts through flow of electrons.

> The n-p boundary has a depletion region.

> If  $V_{GS} < 0$ , as  $V_{GS}$  increases, the depletion region decreases.  $\Rightarrow$  Flow of  $e^-$  from source to drain decreases.  $\Rightarrow$  we get voltage controlled resistor.

Please do not write anything except the question number in this space)

कृपया इस स्थान में प्रश्न संख्या के अनिश्चित कुछ न लिखें।

# UPSC

Answer Questions in NOT MORE THAN the Word Limit specified for each in the Parenthesis.  
Content of the Question is more important than length.  
( Specimen Answer Booklet - For Practice Purpose Only)

उम्मीदवारों को इस हाशिए में नहीं लिखना चाहिए।  
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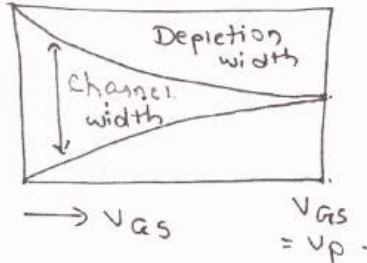


Fig 2. Variation of channel with  $V_{GS}$

> Pinch-off voltage  
 $V_p$  = voltage at which drain current ( $I_D$ ) = 0 i.e. depletion region becomes so large that it completely

closes the n-channel.

> The voltages are related as :-

$$I_D = I_{DSS} \left[ 1 - \frac{V_{GS}}{V_p} \right]^2 \text{ where}$$

$V_{GS}$  = Gate-source voltage

$I_{DSS}$  = Saturation current

$V_p$  = Pinch-off voltage

$I_D$  = Drain current

~~$11 \frac{1}{2}$   
20~~

Advantages of JFET v/s BJT :- ① Higher input impedance  $\Rightarrow$  Low noise ② High power gain  $\Rightarrow$  No need of driver circuit.

Usually, IC 2N547 is used as JFET.

